

# Temperature dependent EXAFS study of $\text{Mn}_5\text{Si}_3\text{C}_x$

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The intermetallic compound  $\text{Mn}_5\text{Si}_3$  exhibits antiferromagnetic order at temperatures below 100 K. The magnetic transition is accompanied with an orthorhombic distortion of the hexagonal unit cell [1]. Further structural and magnetic analysis on single-crystalline samples have shown that a second antiferromagnetic transition occurs at 66 K together with a change of the  $c$  axis dimension [1]. Remarkably, the antiferromagnetic  $\text{Mn}_5\text{Si}_3$  becomes ferromagnetic by doping with carbon [2]. For  $\text{Mn}_5\text{Si}_3\text{C}_x$  films prepared by magnetron sputtering on at substrate temperatures of  $T_S = 740$  K a maximum Curie temperature of  $T_C = 350$  K was found for  $x \approx 1$  and the local magnetic Mn moments were determined by zero-field NMR measurements [3, 4]. Earlier structural investigations performed at room temperature provided strong evidence that the carbon is incorporated in the octahedral voids of the hexagonal  $\text{D}_{8h}$  structure (space group  $\text{P6}_3/\text{mcm}$ ) surrounded by Mn atoms [5]. However, it is not clear whether or not a structural phase transition occurs below room temperature similar to the parent  $\text{Mn}_5\text{Si}_3$  compound. In this respect, conventional x-ray powder diffractometry is an inadequate method to check for a small orthorhombic distortion, since only the peak intensities are slightly changed which are furthermore reduced by the film texture. Therefore, we have started an extended x-ray absorption fine structure (EXAFS) study on  $\text{Mn}_5\text{Si}_3\text{C}_x$  to investigate the local structural properties in detail. A previous EXAFS study performed at ambient temperature suggested a change of the local Mn environment by the incorporation of carbon in the octahedral voids [4]. However, a detailed analysis of the data in terms of a multiple scattering approach by FEFF6 code was not successful and was presumably hampered by the large thermal attenuation of the EXAFS oscillations.

200-nm thick polycrystalline  $\text{Mn}_5\text{Si}_3$  and  $\text{Mn}_5\text{Si}_3\text{C}_{0.75}$  films were prepared by magnetron sputtering on (1-102)-oriented sapphire substrates at a substrate temperature  $T_S = 740$  K. In addition, an  $\text{Mn}_5\text{Si}_3$  powder sample was prepared by sintering a mixed powder of the elemental constituents at 1170 K for eight days for comparison with the  $\text{Mn}_5\text{Si}_3$  film. X-ray absorption spectra at the Mn K-edge were taken at photon energies 6400 – 7500 eV at beamline A1 in the two-crystal mode using a Ge fluorescence detector. Contribution due to the Fe K fluorescence radiation from the background was minimized by electronic discrimination of the Ge detector signal. Multiple EXAFS scans were taken at temperatures  $T = 5$  K, 82 K, and 150 K employing a liquid He flow-cryostat. All spectra were calibrated with an Mn foil (K-edge at 6537.7 eV) and were corrected for self-absorption effects according to Ref. [6] prior to further data processing. The obtained EXAFS oscillations  $\chi(k)k^3$  were Fourier transformed into R space by application of a Bessel window.

Fig. 1 shows the Fourier transforms  $\text{FT}(\chi k^3)$  of the different  $\text{Mn}_5\text{Si}_3\text{C}_x$  samples at various temperatures  $T$ . For  $T = 150$  K all FT show a similar behavior indicating a similar short range order around Mn in the hexagonal  $\text{D}_{8h}$  structure. This is also observed for  $T = 82$  K indicating that the orthorhombic distortion of the crystallographic  $a$  and  $b$  axes reported for single-crystalline samples [1] is very small in the investigated polycrystalline samples. In contrast, the data measured at  $T = 5$  K exhibit clear differences when comparing the FT of the films with the FT of the powder sample, in particular around  $R = 3.7$  Å. However, the qualitative behavior of  $\text{FT}(R)$  of the films seems to be independent of temperature. This demonstrates that the films exhibit do not exhibit an orthorhombic distortion - even not for the undoped  $\text{Mn}_5\text{Si}_3$  film - in contrast to the bulk  $\text{Mn}_5\text{Si}_3$  powder sample. Moreover, the local structural order of  $\text{Mn}_5\text{Si}_3$  films is similar to  $\text{Mn}_5\text{Si}_3\text{C}_{0.75}$  films. For a more detailed analysis the experimental data should be compared with the FT derived from a theoretical calculation of the EXAFS signal by a multiple scattering approach using FEFF6 code [7]. Preliminary calculations have been performed for the films, see Fig. 1. For this purpose we have used the lattice parameters measured by x-ray diffraction at  $T = 25$  K assuming a hexagonal  $\text{D}_{8h}$  structure. Thermal broadening was taken into account by application of the Debye model with a Debye temperature  $\theta_D = 400$  K. The agreement between the experimental data at  $T = 5$  K and the FEFF calculation is very good for the  $\text{Mn}_5\text{Si}_3$  film whereas for the carbon doped  $\text{Mn}_5\text{Si}_3\text{C}_{0.75}$  film differences are found around  $R = 3.8$  Å. Further data analysis is necessary to refine the structural parameters by successive fitting the theoretical FEFF calculations to the experimental data.

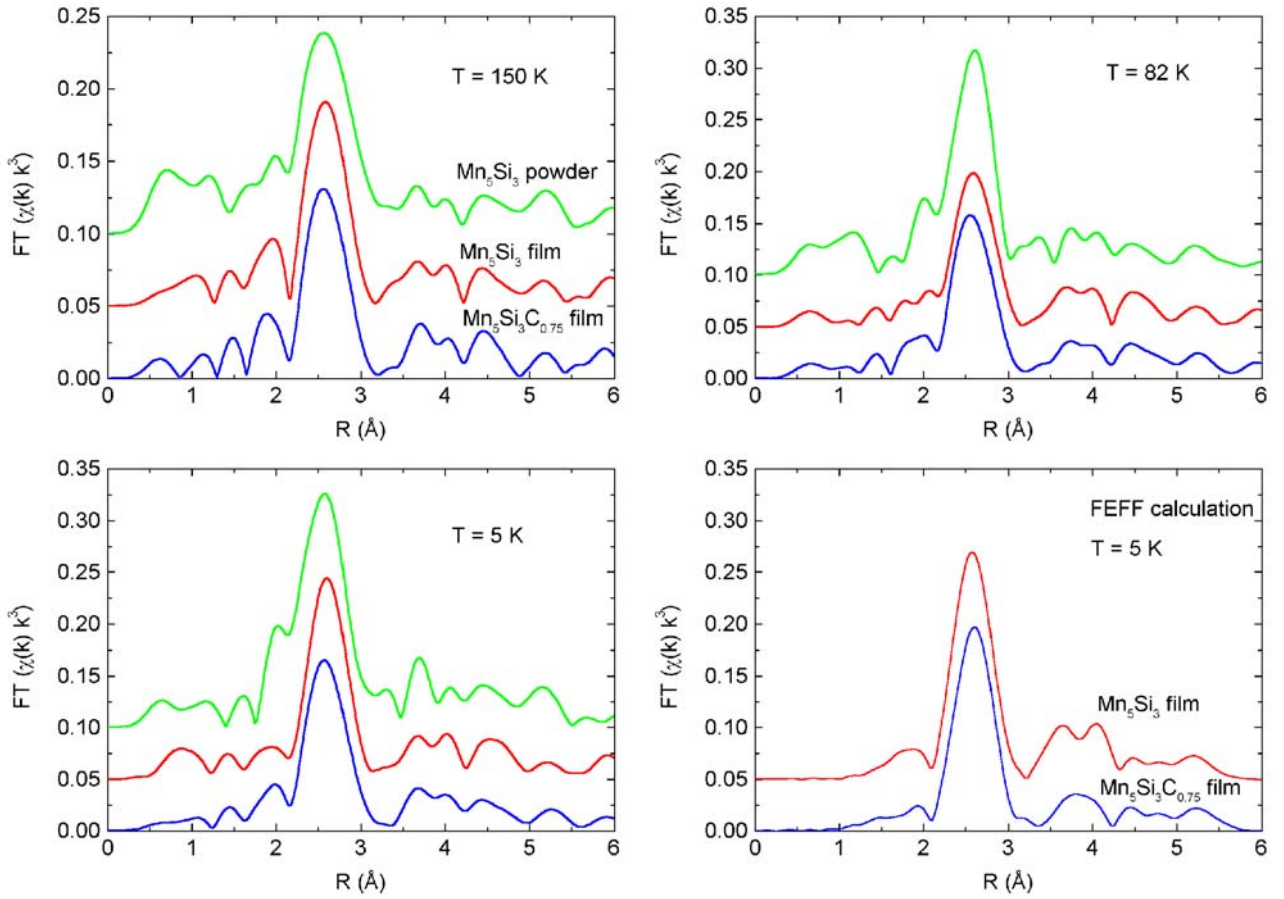


Figure 1: Fourier transforms of the Mn  $K$ -edge EXAFS oscillations  $\chi(k)k^3$  for different  $\text{Mn}_5\text{Si}_3\text{C}_x$  samples at various temperatures  $T$ . In each panel the upper two curves are successively shifted upward by 0.05 for clarity. Also shown are the transforms for  $\text{Mn}_5\text{Si}_3$  and  $\text{Mn}_5\text{Si}_3\text{C}_{0.75}$  at  $T = 5$  K derived from a multiple scattering calculation of the EXAFS signal using FEFF6 code.

## References

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